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La₂LiMoO₆

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Abstract

Dilanthanum lithium molybdenum(V) oxide has a monoclinically distorted perovskite-type structure with an ordered distribution of the Mo⁵⁺ and Li⁺ ions over the octahedral sites. The Mo—O distances within the slightly distorted MoO₆ octahedra range between 1.959 (2) and 1.970 (2) Å with an average value of 1.965 Å. The LiO₆ octahedra are more distorted with Li—O distances in the range 2.099 (2)–2.155 (2) Å (average distance 2.132 Å). The La³⁺ ions are surrounded by eight O atoms at distances ranging from 2.409 (3) to 2.776 (2) Å with the four nearest O atoms forming a distorted tetrahedron and the remaining O atoms roughly capping two faces or bridging two edges of this tetrahedron.

Comment

The AMoO₃ (A = Ca, Sr and Ba) (Scholder & Brixner, 1955) compounds are the only ternary reduced molybde-

num oxides which crystallize with the perovskite structure. Replacement of half of the Mo⁴⁺ ions by a 3d element (Sc, Cr, Fe) (Patterson, Moeller & Ward, 1963) or a rare earth element (La → Lu) (Brandle & Steinfink, 1971) leads to reduced Mo compounds with general formula A₂BMoO₆. The differences in radius and valence between the B³⁺ and Mo⁵⁺ ions lead to the ordering of the latter. As a consequence, these quaternary oxides crystallize in cubic, tetragonal or orthorhombic perovskite supercells, closely related to the cryolite structure, depending upon the sizes of the A and B ions. More complex quadruple oxides containing Mo⁵⁺ ions and having perovskite-related structures are also known, such as ALaMnMoO₆ (A = Ca, Sr and Ba) (Nakamura & Choy, 1977) and La₂M_{4/3}Mo_{2/3}O₆ (M = Mg and Mn) (Bazuev, Makarova & Shveikin, 1987). However, with the exception of Ba₂NdMoO₆, the structure of which was determined at 473 K (Brandle & Steinfink, 1971), none of the above compounds has been the subject of accurate single-crystal or powder structure determinations.

We present here the crystal structure of a new ordered A₂BMoO₆ compound containing Mo⁵⁺ ions: La₂LiMoO₆. This compound belongs to the large family of analogous phases La₂LiMO₆ (M = V, Fe, Nb, Ru, Ta, Re, Os and Ir) and was mentioned by Demazeau, Oh-Kim, Choy & Hagemuller in 1987, though no details of the synthesis or the structural parameters were given. La₂LiMoO₆ crystallizes in a monoclinically distorted variant (space group *P2₁/n*) of the perovskite structure with unit-cell parameters of *ca* $\sqrt{2}a_p \times \sqrt{2}a_p \times 2a_p$ where *a_p* is the unit-cell parameter of an ideal cubic

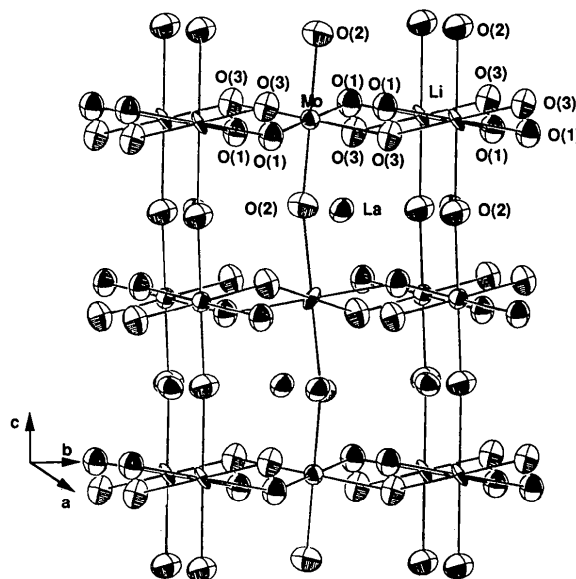


Fig. 1. *ORTEP* (Johnson, 1976) drawing of the structure of La₂LiMoO₆. Displacement ellipsoids are drawn at the 95% probability level.

perovskite. Consequently, the structure consists of a framework of ordered corner-sharing MoO₆ and LiO₆ octahedra which are tilted and rotated to optimize the La—O bond lengths (Fig. 1).

Although they are of low symmetry (point group $\bar{1}$), the MoO₆ octahedra are only slightly distorted with Mo—O distances ranging from 1.959 (2) to 1.970 (2) Å and O—Mo—O angles varying between 88.9 (1) and 89.8 (1)°. The average Mo—O distance is 1.965 Å and is close to that calculated for La₃MoO₇ (1.987 Å) (Tortelier & Gougeon, 1995) in which Mo⁵⁺ ions are also present. On the other hand, this value is smaller than the distance expected from the sum of the ionic radii of O²⁻ and Mo⁵⁺ (2.03 Å) according to Shannon & Prewitt (1969) and larger than the value of 1.86 (2) Å observed for Ba₂NdMoO₆. [The short Mo—O distance in Ba₂NdMoO₆ was attributed to a large overlap of the oxygen orbitals and the 4d orbitals of molybdenum by the authors. However, it could be due to the inaccuracy in the position of the unique crystallographically independent O atom, which shows a large isotropic displacement parameter of 4.1 (4) Å².]

The LiO₆ octahedra are more distorted with Li—O distances ranging between 2.099 (2) and 2.155 (2) Å and O—Li—O angles between 88.7 (1) and 89.92 (9)°. The average Li—O distance is 2.132 Å, which lies between the distances calculated from the sum of the ionic radii of Li⁺ and O²⁻ with coordination number IV (2.12 Å) and VI (2.14 Å).

The deformation from the ideal perovskite structure is mainly due to the small radius of the La³⁺ ion. This is reflected by its environment, which consists of eight O atoms instead of twelve in the ideal cubic perovskite structure. Of these eight O atoms, the nearest four [at distances ranging from 2.409 (3) to 2.496 (3) Å] form a distorted tetrahedron and the remaining four [2.660 (2)–

2.776 (2) Å] roughly cap two faces or bridge two edges of this tetrahedron.

Experimental

Single crystals were obtained by mixing La₂O₃, Li₂MoO₄, MoO₃ and Mo, all in powder form, in the ratio 6:3:2:1. The mixture was pressed into a pellet and placed in a molybdenum crucible which was then sealed under a low argon pressure using an arc welding system. The mixture was heated at a rate of 100 K h⁻¹ to about 2020 K, held at this temperature for 48 h, and then cooled at a rate of 100 K h⁻¹ to 1400 K, at which temperature the power was turned off.

Crystal data

La₂LiMoO₆
M_r = 476.70
 Monoclinic
*P*2₁/*n*
a = 5.590 (1) Å
b = 5.6932 (8) Å
c = 7.878 (2) Å
 β = 90.260 (9)°
V = 250.72 (8) Å³
Z = 2
D_x = 6.314 Mg m⁻³

Mo *K*α radiation
 λ = 0.71073 Å
 Cell parameters from 25 reflections
 θ = 5.2–18.6°
 μ = 19.178 mm⁻¹
T = 295 K
 Irregular
 0.12 × 0.10 × 0.10 mm
 Black

Data collection

Enraf–Nonius CAD-4 diffractometer
 θ –2 θ scans
 Absorption correction: ψ scans (North, Phillips & Mathews, 1968)
T_{min} = 0.678, *T_{max}* = 0.997
 1792 measured reflections
 1519 independent reflections
 1519 observed reflections

R_{int} = 0.033
 θ_{\max} = 40°
h = 0 → 10
k = 0 → 10
l = -14 → 14
 3 standard reflections
 frequency: 90 min
 intensity decay: <1%

Refinement

Refinement on *F*
R = 0.026
wR = 0.033
S = 0.764
 1519 reflections
 50 parameters
 $w = 1/[\sigma^2(F_o) + 0.0004F_o^2 + 1]$ (Killeen & Lawrence, 1969)
 $(\Delta/\sigma)_{\max} < 0.01$

$\Delta\rho_{\max} = 2.50 \text{ e } \text{Å}^{-3}$
 $\Delta\rho_{\min} = -3.11 \text{ e } \text{Å}^{-3}$
 Extinction correction: Stout & Jensen (1968)
 Extinction coefficient: $9.6 (2) \times 10^{-7}$
 Atomic scattering factors from Cromer & Waber (1974)

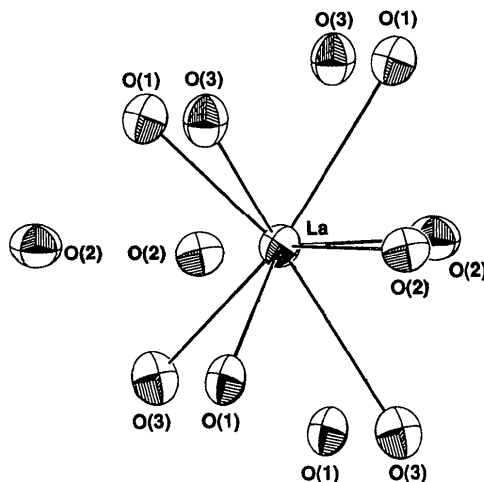


Fig. 2. Environment of the La³⁺ cation. The unconnected ellipsoids represent the remaining four O atoms forming the A site in the ideal cubic perovskite structure.

Table 1. Fractional atomic coordinates and equivalent isotropic displacement parameters (Å²)

	$B_{\text{eq}} = (4/3)\sum_i \sum_j \beta_{ij} a_i \cdot a_j$			
	<i>x</i>	<i>y</i>	<i>z</i>	<i>B_{eq}</i>
Li	0	0	0	0.36 (8)
Mo	1/2	1/2	0	0.422 (4)
La	0.99084 (3)	0.55103 (3)	0.75150 (2)	0.660 (2)
O(1)	0.7139 (4)	0.8069 (4)	0.5420 (3)	0.69 (3)
O(2)	0.9154 (5)	0.0205 (5)	0.2589 (3)	0.82 (3)
O(3)	0.6917 (4)	0.7817 (4)	0.9534 (3)	0.77 (3)

Table 2. Selected geometric parameters (Å, °)

Mo—O(1)	1.970 (2)	O(1)—Mo—O(2)	89.8 (1)
Mo—O(2)	1.959 (2)	O(1)—Mo—O(3)	88.9 (1)
Mo—O(3)	1.965 (2)	O(2)—Mo—O(3)	89.8 (1)
Li—O(1)	2.142 (2)	O(1)—Li—O(2)	88.7 (1)
Li—O(2)	2.099 (2)	O(1)—Li—O(3)	89.92 (9)
Li—O(3)	2.155 (2)	O(2)—Li—O(3)	89.2 (1)
		Mo—O(1)—Li	151.9 (1)
		Mo—O(2)—Li	152.2 (1)
		Mo—O(3)—Li	151.1 (1)
La—O(1)	2.430 (2)	La—O(1)	2.687 (2)
La—O(2)	2.409 (3)	La—O(1)	2.724 (2)
La—O(2)	2.496 (3)	La—O(3)	2.660 (2)
La—O(3)	2.446 (2)	La—O(3)	2.776 (2)

For ease of comparison with the perovskite structure, the non-conventional *P2₁/n* setting was retained. Data were corrected for Lorentz and polarization effects. The structure was solved with the aid of *MULTAN11/82* (Main, Fiske, Hull, Lessinger, Germain, Declercq & Woolfson, 1982) and subsequent difference Fourier syntheses. Refinements of the occupancy factors for the Li and Mo sites confirmed that they are fully occupied. All calculations were performed on a Digital MicroVAX 3100.

Data collection: *CAD-4 Software* (Enraf-Nonius, 1989). Cell refinement: *CAD-4 Software*. Data reduction: *MolEN* (Fair, 1990). Program(s) used to refine structure: *MolEN*. Molecular graphics: *ORTEPII* (Johnson, 1976). Software used to prepare material for publication: *MolEN*.

Lists of structure factors, anisotropic displacement parameters and complete geometry have been deposited with the IUCr (Reference: DU1142). Copies may be obtained through The Managing Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

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Cadmium Hexacyanoiridate(III) Dodecahydrate

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Abstract

Single crystals of triscadmium bis[hexacyanoiridate(III)] dodecahydrate, Cd₃[Ir(CN)₆]₂·12H₂O, were synthesized by the 'slow-diffusion U-tube' method. The disordered structure crystallizes in the cubic space group *Fm $\bar{3}$ m* (No. 225) with a lattice constant of 10.848 (1) Å and one-and-a-third molecules in the unit cell. Divalent Cd atoms are linked to the Ir atoms by linear cyanide bridges. When disorder occurs cavities are formed, a water molecule, O2, replaces an Ir atom and a cyanide bridge is replaced with a coordinated water molecule, O1. A hydrogen-bonding network is set up between O1, O2, and O3 within the cavities of the crystal lattice. Dehydration and IR studies were conducted on single crystals of the title compound. The compound is isostructural with the dodecahydrates of Cd₃[Co(CN)₆]₂ and Mn₃[Co(CN)₆]₂.

Comment

IR, dehydration, electron diffraction and powder and single-crystal X-ray diffraction methods have been employed to investigate the structural and morphological relationships of the transition metal ferro-, ferri- and cobaltcyanides (Milligan *et al.*, 1970; Beall, Milligan, Korp & Bernal, 1977; Mullica, Oliver, Milligan & Hills, 1979). However, in spite of the structural investigations of the hydrated Prussian blue analogs, there are still questions concerning the zeolitic water structure in these materials. A clear understanding of the water network is essential in order to model the H₂O transport in the semi-permeable membranes. Several structural investigations of the transition metal hexacyanides have led to proposals of various models for the face-centered cubic lattice (Keggin & Miles, 1936; Ludi & Güdel,